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About a regular point of the lattice spin-boson model with at most two photons

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The spin-boson model is a well-known quantum-mechanical model which describes the interaction between a two-level atom and a photon field. We refer to [1] and [2] for excellent reviews respectively from physical and mathematical perspectives.

Let T^d be the d -dimensional torus, $L_2((T^d))$ be the Hilbert space of square integrable (complex) functions defined on T^d , \mathbb{C}^2 be the state of the two-level atom and $\mathcal{F}_b(L_2(T^d))$ be the symmetric Fock space for bosons, that is,

$$\mathcal{F}_b(L_2(T^d)) := \mathbb{C} \oplus L_2(T^d) \oplus L_2^{sym}((T^d)^2) \oplus \dots$$

Here $L_2^{sym}((T^d)^n)$ is the Hilbert space of symmetric functions of $n \geq 2$ variables. For $m = 1, 2$ we denote $\mathcal{L}_m := \mathbb{C}^2 \otimes \mathcal{F}_b^{(m)}(L_2(T^d))$, where

$$\mathcal{F}_b^{(1)}(L_2(T^d)) := \mathbb{C} \oplus L_2(T^d), \quad \mathcal{F}_b^{(2)}(L_2(T^d)) := \mathbb{C} \oplus L_2(T^d) \oplus L_2^{sym}((T^d)^2).$$

We write elements F of the space \mathcal{L}_2 in the form $F = \{f_0^{(s)}, f_1^{(s)}(k_1), f_2^{(s)}(k_1, k_2); s = \pm\}$. Then the norm in \mathcal{L}_2 is given by

$$\|F\|^2 := \sum_{s=\pm} \left(|f_0^{(s)}|^2 + \int_{T^d} |f_1^{(s)}(k_1)|^2 dk_1 + \frac{1}{2} \int_{(T^d)^2} |f_2^{(s)}(k_1, k_2)|^2 dk_1 dk_2 \right). \quad (1)$$

We recall that the lattice spin-boson model with at most two photons \mathcal{A}_2 is acting in \mathcal{L}_2 as the 3×3 tridiagonal block operator matrix

$$\mathcal{A}_2 := \begin{pmatrix} A_{00} & A_{01} & 0 \\ A_{01}^* & A_{11} & A_{12} \\ 0 & A_{12}^* & A_{22} \end{pmatrix},$$

where matrix elements A_{ij} are defined by

$$\begin{aligned}
A_{00}f_0^{(s)} &= s\varepsilon f_0^{(s)}, \quad A_{01}f_1^{(s)} = \alpha \int_{T^d} v(t)f_1^{(-s)}(t)dt, \\
(A_{11}f_1^{(s)})(k_1) &= (s\varepsilon + w(k_1))f_1^{(s)}(k_1), \quad (A_{12}f_2^{(s)})(k_1) = \alpha \int_{T^d} v(t)f_2^{(-s)}(k_1, t)dt, \\
(A_{22}f_2^{(s)})(k_1, k_2) &= (s\varepsilon + w(k_1) + w(k_2))f_2^{(s)}(k_1, k_2), \\
f &= \{f_0^{(s)}, f_1^{(s)}, f_2^{(s)}; s = \pm\} \in \mathcal{L}_2.
\end{aligned}$$

Here A_{ij}^* denotes the adjoint operator to A_{ij} for $i < j$ with $i, j = 0, 1, 2$; $w(k)$ is the dispersion of the free field, $\alpha v(k)$ is the coupling between the atoms and the field modes, $\alpha > 0$ is a real number, so-called the coupling constant, real number. We assume that $v(\cdot)$ and $w(\cdot)$ are the real-valued continuous functions on T^d . Under these assumptions the lattice spin-boson model with at most two photons \mathcal{A}_2 is bounded and self-adjoint in the complex Hilbert space \mathcal{L}_2 .

To study the spectral properties of \mathcal{A}_2 we introduce the following two bounded self-adjoint operators $\hat{\mathcal{A}}_2^{(s)}$, $s = \pm$, which acts in $\mathcal{F}_b^{(2)}(L_2(T^d))$ as

$$\hat{\mathcal{A}}_2^{(s)} := \begin{pmatrix} \hat{A}_{00}^{(s)} & \hat{A}_{01} & 0 \\ \hat{A}_{01}^* & \hat{A}_{11}^{(s)} & \hat{A}_{12} \\ 0 & \hat{A}_{12}^* & \hat{A}_{22}^{(s)} \end{pmatrix}$$

with the entries

$$\begin{aligned}
\hat{A}_{00}^{(s)}f_0 &= s\varepsilon f_0, \quad \hat{A}_{01}f_1 = \alpha \int_{T^d} v(t)f_1(t)dt, \\
(\hat{A}_{11}^{(s)}f_1)(k_1) &= (-s\varepsilon + w(k_1))f_1(k_1), \quad (\hat{A}_{12}f_2)(k_1) = \alpha \int_{T^d} v(t)f_2(k_1, t)dt, \\
(\hat{A}_{22}^{(s)}f_2)(k_1, k_2) &= (s\varepsilon + w(k_1) + w(k_2))f_2(k_1, k_2), \quad (f_0, f_1, f_2) \in \mathcal{F}_b^{(2)}(L_2(T^d)).
\end{aligned}$$

It is easy to check that

$$\begin{aligned}
(\hat{A}_{01}^*f_0)(k_1) &= \alpha v(k_1)f_0; \\
(\hat{A}_{12}^*f_1)(k_1, k_2) &= \alpha(v(k_1)f_1(k_2) + v(k_2)f_1(k_1)), \quad (f_0, f_1) \in \mathcal{F}_b^{(1)}(L_2(T^d)).
\end{aligned}$$

In order to describe the essential spectrum of \mathcal{A}_2 we define an analytic function $\Delta^{(s)}(\cdot)$ in $\mathbb{C} \setminus [s\varepsilon + m; s\varepsilon + M]$ by

$$\Delta^{(s)}(\lambda) := -s\varepsilon - \lambda - \alpha^2 \int_{T^d} \frac{v^2(t)dt}{s\varepsilon + w(t) - \lambda},$$

where the numbers m and M are defined by

$$m := \min_{p \in T^d} w(p), \quad M := \max_{p \in T^d} w(p).$$

Let $\sigma^{(s)}$ be the set of all complex numbers $\lambda \in \mathbb{C}$ such that the equality $\Delta^{(s)}(\lambda - w(k_1)) = 0$ holds for some $k_1 \in T^d$. Then for the essential spectrum of $\mathcal{A}_2^{(s)}$ we have

$$\sigma_{ess}(\mathcal{A}_2^{(s)}) = \sigma^{(s)} \cup [s\varepsilon + 2m; s\varepsilon + 2M].$$

We set $E_{\min}^{(s)} := \min \sigma_{ess}(\mathcal{A}_2^{(s)})$, for $s = \pm$ and $E_{\min} := \min\{E_{\min}^{(+)}, E_{\min}^{(-)}\}$.

Next, we represent the space $\mathcal{F}_b^{(2)}(L_2(T^d))$ as a direct sum of two Hilbert spaces $\mathcal{F}_b^{(1)}(L_2(T^d))$ and $L_2^{sym}((T^d)^2)$, that is, $\mathcal{F}_b^{(2)}(L_2(T^d)) = \mathcal{F}_b^{(1)}(L_2(T^d)) \oplus L_2^{sym}((T^d)^2)$. Then the first Schur complement of the operator $\mathcal{A}_2^{(s)}$ with respect to this decomposition is defined as

$$S_1^{(s)}(\lambda) : \mathcal{F}_b^{(1)}(L_2(T^d)) \rightarrow \mathcal{F}_b^{(1)}(L_2(T^d)), \quad \lambda \in \rho(\hat{A}_{22}^{(s)});$$

$$S_1^{(s)}(\lambda) := \begin{pmatrix} \hat{A}_{00}^{(s)} & \hat{A}_{01} \\ \hat{A}_{01}^* & \hat{A}_{11}^{(s)} \end{pmatrix} - \lambda - \begin{pmatrix} 0 \\ \hat{A}_{12} \end{pmatrix} (\hat{A}_{22}^{(s)} - \lambda)^{-1} \begin{pmatrix} 0 & \hat{A}_{12}^* \end{pmatrix}.$$

Define

$$S_{00}^{(s)}(\lambda) := \hat{A}_{00}^{(s)} - \lambda, \quad S_{01}^{(s)}(\lambda) := \hat{A}_{01};$$

$$S_{10}^{(s)}(\lambda) := \hat{A}_{01}^*, \quad S_{11}^{(s)}(\lambda) := \hat{A}_{11}^{(s)} - \lambda - \hat{A}_{12}(\hat{A}_{22}^{(s)} - \lambda)^{-1} \hat{A}_{12}^*.$$

Then the operator $S_1^{(s)}(\lambda)$ has form

$$S_1^{(s)}(\lambda) = \begin{pmatrix} S_{00}^{(s)}(\lambda) & S_{01}^{(s)}(\lambda) \\ S_{10}^{(s)}(\lambda) & S_{11}^{(s)}(\lambda) \end{pmatrix}.$$

We study some important properties of the first Schur complement

$$S_1(\lambda) := \text{diag}\{S_1^{(+)}(\lambda), S_1^{(-)}(\lambda)\}$$

for the lattice spin-boson model with at most two photons \mathcal{A}_2 .

Proposition 1. *The number $\lambda \in \mathbb{C} \setminus \sigma_{ess}(\mathcal{A}_2)$ is an eigenvalue of the operator \mathcal{A}_2 if and only if the operator $S_1(\lambda)$ has an eigenvalue equal to zero. Moreover, the eigenvalues λ and 0 have the same multiplicities.*

Proposition 2. $\lambda \in \sigma_{ess}(\mathcal{A}_2) \setminus \sigma(A_{22})$ if and only if $0 \in \sigma_{ess}(S_1(\lambda))$.

From Propositions 1 and 2 we obtain the following two corollaries.

Corollary 1. *Let $\lambda \in \mathbb{C} \setminus \sigma_{ess}(\mathcal{A}_2)$. Then $\lambda \in \rho(\mathcal{A}_2) \Leftrightarrow 0 \in \rho(S_1(\lambda))$.*

Corollary 2. *Let $\lambda_0 \in \mathbb{R} \setminus \sigma_{ess}(\mathcal{A}_2)$. If $(\lambda_0; \lambda_0 + \gamma) \subset \rho(\mathcal{A}_2)$ (resp. $(\lambda_0 - \gamma; \lambda_0) \subset \rho(\mathcal{A}_2)$) for some $\gamma > 0$, then there exists a number $\delta = \delta(\gamma) > 0$ such that $(0; \delta) \subset \rho(S_1(\lambda_0))$ (resp. $(-\delta; 0) \subset \rho(S_1(\lambda_0))$).*

For a bounded self-adjoint operator \mathcal{A} acting in a Hilbert space H we denote by $N_{(-\infty; \lambda)}(\mathcal{A})$ the number of eigenvalues of \mathcal{A} to the left of λ , $\lambda \leq \min \sigma_{ess}(\mathcal{A})$.

Proposition 3. *The function $N_{(-\infty; 0)}(S_1(\cdot))$ is monotonically increasing in $(-\infty; 0)$.*

Definition 1. *We denote by $E_m(\cdot), m \in \mathbb{N}$ the positive definite function on the segment $[\alpha; \beta] \subset \mathbb{R} \setminus \sigma(\mathcal{A}_2)$, satisfying the condition: $E_m(\lambda)$ is the m -th eigenvalue (eigenvalues numbered in ascending order, counting their multiplicity) of the operator $S_1(\lambda), \lambda \in [\alpha; \beta]$.*

Theorem. *The number $\lambda_0 < E_{\min}$ is the regular point of the operator \mathcal{A}_2 if and only if the function $N_{(-\infty; 0)}(S_1(\cdot))$ is continuous at point $\lambda = \lambda_0$.*

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