



**FIZIKA, MATEMATIKA VA  
MEXANIKANING DOLZARB  
MUAMMOLARI  
XALQARO ILMIY-AMALIY  
ANJUMANI  
MATERIALLARI**

**BUXORO DAVLAT UNIVERSITETI**

**Buxoro - 2023**

**O‘ZBEKISTON RESPUBLIKASI OLIIY TA‘LIM, FAN VA  
INNOVATSIYALAR VAZIRLIGI  
BUXORO DAVLAT UNIVERSITETI**

**FIZIKA, MATEMATIKA VA MEKANIKA DOLZARB  
MUAMMOLARI**

xalqaro ilmiy-amaliy anjumani

## **MATERIALLARI**

**(I qism)**

Buxoro, O‘zbekiston, 24-25-may, 2023-yil

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**МИНИСТЕРСТВО ВЫСШЕГО ОБРАЗОВАНИЯ, НАУКИ И  
ИННОВАЦИЙ РЕСПУБЛИКИ УЗБЕКИСТАН  
БУХАРСКИЙ ГОСУДАРСТВЕННЫЙ УНИВЕРСИТЕТ**

## **ТЕЗИСЫ ДОКЛАДОВ**

**(Часть I)**

международной научно-практической конференции

**АКТУАЛЬНЫЕ ПРОБЛЕМЫ ФИЗИКИ, МАТЕМАТИКИ И  
МЕХАНИКИ**

Бухара, Узбекистан, 24-25 мая, 2023 год

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**MINISTRY OF HIGHER EDUCATION, SCIENCE AND INNOVATIONS  
OF THE REPUBLIC OF UZBEKISTAN  
BUKHARA STATE UNIVERSITY**

## **ABSTRACTS**

**(Part I)**

of the international scientific and practical conference

**ACTUAL PROBLEMS OF PHYSICS, MATHEMATICS AND  
MECHANICS**

Bukhara, Uzbekistan, May 24-25, 2023

**Fizika, matematika va mexanikaning dolzarb muammolari** (Xalqaro ilmiy-amaliy konfferensiya materiallari to‘plami, I qism) Buxoro-2023, 357 bet.

Mazkur to‘plam “Fizika, matematika va mexanikaning dolzarb muammolari” Xalqaro ilmiy-amaliy konferensiyasi materiallari to‘plami asosida tayyorlangan bo‘lib, matematik analiz, differensial tenglamalar va matematik fizika, algebra va geometriya, hisoblash matematikasi va mexanika, geofizika va qayta tiklanuvchi energiya manbalari, kondensirlangan holatlar fizikasi, zamonaviy ta’limda raqamli texnologiyalar, ehtimollar nazariyasi va matematik statistika yo‘nalishlaridagi ilmiy ma’ruzalar o‘rin olgan.

*To‘plamga kiritilgan maqola va tezislar mazmuni, ilmiyligi va dalillarining haqqoniyligi uchun mualliflar ma’suldirlar*

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**LOCATION OF THE ESSENTIAL SPECTRUM OF THE CHANNEL  
OPERATOR CORRESPONDING TO THE THIRD-ORDER OPERATOR  
MATRIX IN THE FERMIONIC FOCK SPACE**

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Denote by  $\mathbb{C}$  the one-dimensional complex space, by  $L_2(\mathbb{T}^d)$  the Hilbert space of square integrable (complex) functions defined on  $\mathbb{T}^d$  and by  $L_2^{as}((\mathbb{T}^d)^2)$  the Hilbert space of antisymmetric square integrable (complex) functions defined on  $(\mathbb{T}^d)^2$ , that is,  $f(p, q) = -f(q, p)$  for all  $p, q \in \mathbb{T}^d$

For convenience, we introduce the following spaces:

$$\mathcal{F}_a^{(1)}(L_2(\mathbb{T}^d)) := \mathbb{C} \oplus L_2(\mathbb{T}^d);$$

$$\mathcal{F}_a^{(2)}(L_2(\mathbb{T}^d)) := \mathbb{C} \oplus L_2(\mathbb{T}^d) \oplus L_2^{as}((\mathbb{T}^d)^2);$$

$$\mathcal{F}_a(L_2(\mathbb{T}^d)) := \mathbb{C} \oplus L_2(\mathbb{T}^d) \oplus L_2^{as}((\mathbb{T}^d)^2) \oplus \dots$$

**Definition 1.**  $\mathcal{F}_a(L_2(\mathbb{T}^d))$  is called a fermionic Fock space over  $L_2(\mathbb{T}^d)$ ,  $\mathcal{F}_a^{(m)}(L_2(\mathbb{T}^d))$ ,  $m = 1, 2$  is  $m+1$  -particle cut subspace of fermionic Fock space ,  $\mathcal{H}_n$ ,  $n = 0, 1, 2$  space is called the  $n$ -particle subspace of a fermionic Fock space.

Let

$$\mathcal{H}_0 := \mathbb{C}, \mathcal{H}_1 := L_2(\mathbb{T}^d), \mathcal{H}_2 := L_2^{as}((\mathbb{T}^d)^2).$$

We consider the third-order operator matrix defined as

$$\mathcal{A}_{\mu,\lambda} := \begin{pmatrix} A_{00} & \mu A_{01} & 0 \\ \mu A_{01}^* & A_{11} & \mu A_{12} \\ 0 & \mu A_{12}^* & A_{22}^0 - \lambda V \end{pmatrix} \quad (1)$$

in the Hilbert space  $\mathcal{F}_a^{(2)}(L_2(\mathbb{T}^d))$ . Here the matrix elements  $A_{ij}: \mathcal{H}_j \rightarrow \mathcal{H}_i, i, j = 0, 1, 2$  are defined by:

$$A_{00}f_0 = \omega_0 f_0, \quad A_{01}f_1 = \int_{\mathbb{T}^d} v(t) f_1(t) dt,$$

$$(A_{11}f_1)(p) = \omega_1(p) f_1(p), \quad (A_{12}f_2)(p) = \int_{\mathbb{T}^d} v(t) f_2(p, t) dt,$$

$$(A_{22}^0 f_2)(p, q) = \omega_2(p, q) f_2(p, q), \quad V := V_1 + V_2,$$

$$(V_1 f_2)(p, q) = \varphi(q) \int_{\mathbb{T}^d} \varphi(t) f_2(p, t) dt, \quad (V_2 f_2)(p, q) = \varphi(p) \int_{\mathbb{T}^d} \varphi(t) f_2(t, q) dt.$$

Here  $\mu, \lambda$  are real positive numbers,  $\omega_0$  is a fixed real number, functions  $\omega_1(\cdot), v(\cdot)$  and  $\varphi(\cdot)$  are real-valued continuous functions on  $\mathbb{T}^d$ , and  $\omega_2(\cdot, \cdot)$  is a real-valued symmetric continuous function on  $(\mathbb{T}^d)^2$ .

Operator matrix  $\mathcal{A}_{\mu,\lambda}, \mu, \lambda > 0$  acting by formula (1) in the Hilbert space  $\mathcal{F}_a^{(2)}(L_2(\mathbb{T}^d))$  is linear, bounded and self-adjoint operator.

We introduce the following designations:

$$\widehat{\mathcal{H}}_1 := \mathcal{H}_1, \quad \widehat{\mathcal{H}}_2 := L_2((\mathbb{T}^d)^2), \quad \widehat{\mathcal{H}} := \widehat{\mathcal{H}}_1 \oplus \widehat{\mathcal{H}}_2.$$

We consider the second-order operator matrix defined as

$$\mathcal{A}_{\mu,\lambda}^{\text{ch}} := \begin{pmatrix} A_{11} & (\mu/\sqrt{2})A_{12} \\ (\mu/\sqrt{2})A_{12}^* & A_{22}^0 - \lambda V_1 \end{pmatrix}$$

in the Hilbert space  $\widehat{\mathcal{H}}$  and called channel operator, which plays an important role in the analysis of the spectrum of the operator matrix  $\mathcal{A}_{\mu,\lambda}$ .

It can be easily verified that the channel operator  $\mathcal{A}_{\mu,\lambda}^{\text{ch}}$  is linear, bounded and self-adjoint in the Hilbert space  $\widehat{\mathcal{H}}$ .

We consider the generalized Friedrichs model [1]-[3]  $\mathcal{A}_{\mu,\lambda}(k), k \in \mathbb{T}^d$  acting in the space  $\mathcal{H}_0 \oplus \mathcal{H}_1$  as

$$\mathcal{A}_{\mu,\lambda}(k) := \begin{pmatrix} A_{00}(k) & \mu A_{01} \\ \mu A_{01}^* & A_{11}^0(k) - \lambda v \end{pmatrix}.$$

Here

$$\begin{aligned} A_{00}(k)f_0 &= \omega_1(k)f_0, & (A_{11}^0(k)f_1)(q) &= \omega_2(k, q)f_1(q), \\ (vf_1)(q) &= \varphi(q) \int_{\mathbb{T}^d} \varphi(t)f_1(t)dt. \end{aligned}$$

For any  $k \in \mathbb{T}^d$  we define an analytic function

$$\begin{aligned} \Delta_{\mu,\lambda}(k; z) &= \left( 1 - \lambda \int_{\mathbb{T}^d} \frac{\varphi^2(t)dt}{\omega_2(k, t) - z} \right) \left( \omega_1(k) - z - \frac{\mu^2}{2} \int_{\mathbb{T}^d} \frac{v^2(t)dt}{\omega_2(k, t) - z} \right) \\ &\quad - \frac{\lambda\mu^2}{2\sqrt{2}} \left( \int_{\mathbb{T}^d} \frac{v(t)\varphi(t)dt}{\omega_2(k, t) - z} \right)^2 \end{aligned}$$

in  $\mathbb{C} \setminus \sigma_{\text{ess}}(\mathcal{A}_{\mu,\lambda}(k))$ . The function  $\Delta_{\mu,\lambda}(k; z)$  usually called the Fredholm determinant corresponding to the operator matrix  $\mathcal{A}_{\mu,\lambda}(k)$ .

**Lemma 1.** For any  $\mu, \lambda > 0$  and  $k \in \mathbb{T}^d$  the operator  $\mathcal{A}_{\mu,\lambda}(k)$  has an eigenvalue  $z_{\mu,\lambda}(k) \in \mathbb{C} \setminus \sigma_{\text{ess}}(\mathcal{A}_{\mu,\lambda}(k))$  if and only if  $\Delta_{\mu,\lambda}(k, z) = 0$ .

From Lemma 1 it follows that for the discrete spectrum of  $\mathcal{A}_{\mu,\lambda}(k)$  the equality

$$\sigma_{\text{disc}}(\mathcal{A}_{\mu,\lambda}(k)) = \{z \in \mathbb{C} \setminus \sigma_{\text{ess}}(\mathcal{A}_{\mu,\lambda}(k)) : \Delta_{\mu,\lambda}(k, z) = 0\}, k \in \mathbb{T}^d$$

holds.

We introduce the following notations:

$$\sigma_{\text{two}}(\mathcal{A}_{\mu,\lambda}^{\text{ch}}) = \bigcup_{k \in \mathbb{T}^d} \sigma_{\text{disc}}(\mathcal{A}_{\mu,\lambda}(k)), \quad \sigma_{\text{three}}(\mathcal{A}_{\mu,\lambda}^{\text{ch}}) = [m; M],$$

where

$$m := \min_{k, q \in \mathbb{T}^d} \omega_2(k, q), \quad M = \max_{k, q \in \mathbb{T}^d} \omega_2(k, q).$$



**Theorem 1.** The spectrum  $\sigma(\mathcal{A}_{\mu,\lambda}^{\text{ch}})$  of the channel operator  $\mathcal{A}_{\mu,\lambda}^{\text{ch}}$  satisfies the following equality

$$\sigma(\mathcal{A}_{\mu,\lambda}^{\text{ch}}) = \sigma_{\text{two}}(\mathcal{A}_{\mu,\lambda}^{\text{ch}}) \cup \sigma_{\text{three}}(\mathcal{A}_{\mu,\lambda}^{\text{ch}}).$$

**Theorem 2.** The essential spectrum  $\sigma_{\text{ess}}(\mathcal{A}_{\mu,\lambda})$  of the operator matrix  $\mathcal{A}_{\mu,\lambda}$  coincides with the spectrum  $\sigma(\mathcal{A}_{\mu,\lambda}^{\text{ch}})$  of the channel operator matrix  $\mathcal{A}_{\mu,\lambda}^{\text{ch}}$ , i.e.  $\sigma_{\text{ess}}(\mathcal{A}_{\mu,\lambda}) = \sigma(\mathcal{A}_{\mu,\lambda}^{\text{ch}})$ .

**Definition 2.** The sets  $\sigma_{\text{two}}(\mathcal{A}_{\mu,\lambda}^{\text{ch}})$  and  $\sigma_{\text{three}}(\mathcal{A}_{\mu,\lambda}^{\text{ch}})$  are called two- and three-particle branches of the essential spectrum of the operator matrix  $\mathcal{A}_{\mu,\lambda}$ , respectively.

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## ON TRAJECTORIES OF A NON-VOLTERRA QUADRATIC STOCHASTIC OPERATOR

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Let

$$S^{m-1} = \left\{ x = (x_1, x_2, \dots, x_m) \in R^m : x_i \geq 0, \text{ for any } i \text{ and } \sum_{i=1}^m x_i = 1 \right\}$$

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