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(Dilmurodov E.B., Rasulov T.H.) p.7

FORMATION AND DEVELOPMENT
OF COMPETITIVE SKILLS
IN THE SUBJECTS
OF "MASS CULTURE"
IN CONTINUOUS EDUCATIONAL
PROCESS
(Tojiboyeva H.M) p.51

PROFESSIONAL ORIENTATION
OF COMMUNICATIVE
COMPETENCE OF STUDENTS
(Kasimova Z.Kh.) p.53



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EUROPEAN SCIENCE № 2(51) Part II 2020 ISSN 2541-786X

Содержание

| | |
|--|-----------|
| PHYSICO-MATHEMATICAL SCIENCES | 7 |
| <i>Dilmurodov E.B., Rasulov T.H.</i> (Republic of Uzbekistan) ESSENTIAL SPECTRUM OF A 2x2 OPERATOR MATRIX AND THE FADDEEV EQUATION / <i>Дилмуродов Э.Б., Расулов Т.Х.</i> (Республика Узбекистан) СУЩЕСТВЕННЫЙ СПЕКТР ОДНОЙ 2x2 ОПЕРАТОРНОЙ МАТРИЦЫ И УРАВНЕНИЕ ФАДДЕЕВА | 7 |
| <i>Tosheva N.A., Rasulov T.H.</i> (Republic of Uzbekistan) MAIN PROPERTY OF REGULARIZED FREDHOLM DETERMINANT CORRESPONDING TO A FAMILY OF 3x3 OPERATOR MATRICES / <i>Тошева Н.А., Расулов Т.Х.</i> (Республика Узбекистан) ОСНОВНЫЕ СВОЙСТВА РЕГУЛЯРИЗОВАННОГО ОПРЕДЕЛИТЕЛЯ ФРЕДГОЛЬМА, СООТВЕТСТВУЮЩИЕ СЕМЕЙСТВУ 3x3 ОПЕРАТОРНЫХ МАТРИЦ | 11 |
| <i>Bahronov B.I., Rasulov T.H.</i> (Republic of Uzbekistan) STRUCTURE OF THE NUMERICAL RANGE OF FRIEDRICH'S MODEL WITH RANK TWO PERTURBATION / <i>Бахронов Б.И., Расулов Т.Х.</i> (Республика Узбекистан) СТРУКТУРА ЧИСЛОВОЙ ОБЛАСТИ ЗНАЧЕНИЙ МОДЕЛИ ФРИДРИХСА С ДВУМЕРНЫМ ВОЗМУЩЕНИЕМ | 15 |
| <i>Umirkulova G.H., Rasulov T.H.</i> (Republic of Uzbekistan) CHARACTERISTIC PROPERTY OF THE FADDEEV EQUATION FOR THREE-PARTICLE MODEL OPERATOR ON A ONE-DIMENSIONAL LATTICE / <i>Умиркулова Г.Х., Расулов Т.Х.</i> (Республика Узбекистан) ХАРАКТЕРИСТИЧЕСКИЕ СВОЙСТВА УРАВНЕНИЯ ФАДДЕЕВА ДЛЯ ТРЕХЧАСТИЧНОГО МОДЕЛЬНОГО ОПЕРАТОРА НА ОДНОМЕРНОЙ РЕШЕТКЕ | 19 |
| <i>Mustafoeva Z.E., Rasulov T.H.</i> (Republic of Uzbekistan) INVESTIGATION OF THE SPECTRUM OF A DIAGONALIZABLE 4x4-OPERATOR MATRIX / <i>Мустафоева З.Э., Расулов Т.Х.</i> (Республика Узбекистан) ИССЛЕДОВАНИЕ СПЕКТРА ОДНОЙ ДИАГОНАЛИЗИРУЕМОЙ 4x4-ОПЕРАТОРНОЙ МАТРИЦЫ | 23 |
| <i>Merajov N.I., Rasulov T.H.</i> (Republic of Uzbekistan) DESCRIPTION OF THE POINT SPECTRUM OF A 3x3 TRIDIAGONAL OPERATOR MATRIX WITH FREDHOLM OPERATORS / <i>Меражов Н.И., Расулов Т.Х.</i> (Республика Узбекистан) ОПИСАНИЕ ТОЧЕЧНОГО СПЕКТРА ТРИДАГОНАЛЬНОГО 3x3 ОПЕРАТОРНОЙ МАТРИЦЫ С ФРЕДГОЛЬМСКИМИ ОПЕРАТОРАМИ | 27 |
| <i>Nematova Sh.B., Rasulov T.H.</i> (Republic of Uzbekistan) THRESHOLD EIGENVALUES OF A TWO-CHANNEL MOLECULAR-RESONANCE MODEL / <i>Неъматова Ш.Б., Расулов Т.Х.</i> (Республика Узбекистан) ПОРОГОВЫЕ СОБСТВЕННЫЕ ЗНАЧЕНИЯ ДВУХКАНАЛЬНОЙ МОЛЕКУЛЯРНО-РЕЗОНАНСНОЙ МОДЕЛИ | 31 |
| TECHNICAL SCIENCES | 35 |
| <i>Mansurova Sh.P.</i> (Republic of Uzbekistan) QUESTIONS FEATURES OF DESIGNING AIR CURTAIN / <i>Мансурова Ш.П.</i> (Республика Узбекистан) ВОПРОСЫ ОСОБЕННОСТИ ПРОЕКТИРОВАНИЯ ВОЗДУШНЫХ ЗАВЕС | 35 |

| | |
|--|-----------|
| <i>Ustemirov Sh.R.</i> (Republic of Uzbekistan) ANALYSIS OF REVERSE WATER SUPPLY SYSTEMS AND PROBLEMS OF WATER QUALITY OF INDUSTRIAL ENTERPRISES / <i>Устемиров Ш.Р.</i> (Республика Узбекистан) АНАЛИЗ СИСТЕМ ОБОРОТНОГО ВОДОСНАБЖЕНИЯ И ПРОБЛЕМ КАЧЕСТВА ВОДЫ ПРОМЫШЛЕННЫХ ПРЕДПРИЯТИЙ..... | 39 |
| AGRICULTURAL SCIENCES..... | 42 |
| <i>Isaeva L.B., Sanoev H.A.</i> (Republic of Uzbekistan) DYNAMICS OF SOIL HUMIDITY IN THE ROOT TREE OF A PLANT / <i>Исаева Л.Б., Саноев Х.А.</i> (Республика Узбекистан) ДИНАМИКА ВЛАЖНОСТИ ПОЧВЫ В КОРНЕВОМ СТВОЛЕ РАСТЕНИЯ..... | 42 |
| ECONOMICS | 45 |
| <i>Makarenko V.V., Zaporozhtseva E.N.</i> (Russian Federation) FINANCIAL STATEMENTS AS THE MAIN SOURCE OF INFORMATION ON THE FINANCIAL POSITION OF THE ENTERPRISE / <i>Макаренко В.В., Запорожцева Е.Н.</i> (Российская Федерация) БУХГАЛТЕРСКАЯ ОТЧЁТНОСТЬ КАК ОСНОВНОЙ ИСТОЧНИК ИНФОРМАЦИИ О ФИНАНСОВОМ ПОЛОЖЕНИИ ПРЕДПРИЯТИЯ..... | 45 |
| PHILOLOGICAL SCIENCES..... | 49 |
| <i>Karimov Z.A.</i> (Republic of Uzbekistan) PHILOSOPHICAL ANALYSIS OF LIFESTYLE AND REPRODUCTIVE NOTIONS / <i>Каримов З.А.</i> (Республика Узбекистан) ФИЛОСОФСКИЙ АНАЛИЗ ОБРАЗА ЖИЗНИ И РЕПРОДУКТИВНЫХ ПОНЯТИЙ..... | 49 |
| PEDAGOGICAL SCIENCES..... | 51 |
| <i>Tojiboyeva H.M.</i> (Republic of Uzbekistan) FORMATION AND DEVELOPMENT OF COMPETITIVE SKILLS IN THE SUBJECTS OF “MASS CULTURE” IN CONTINUOUS EDUCATIONAL PROCESS / <i>Тожибоева Х.М.</i> (Республика Узбекистан) ФОРМИРОВАНИЕ И РАЗВИТИЕ КОНКУРЕНТНЫХ НАВЫКОВ В СУБЪЕКТАХ «МАССОВОЙ КУЛЬТУРЫ» В НЕПРЕРЫВНОМ ОБРАЗОВАТЕЛЬНОМ ПРОЦЕССЕ | 51 |
| <i>Kasimova Z.Kh.</i> (Republic of Uzbekistan) PROFESSIONAL ORIENTATION OF COMMUNICATIVE COMPETENCE OF STUDENTS / <i>Касимова З.Х.</i> (Республика Узбекистан) ПРОФЕССИОНАЛЬНАЯ НАПРАВЛЕННОСТЬ КОММУНИКАТИВНОЙ КОМПЕТЕНТНОСТИ СТУДЕНТОВ | 53 |
| <i>Kakhkhorov S.K., Mirzoyev D.P.</i> (Republic of Uzbekistan) RESEARCHING COMMUTATION DEVICES / <i>Каххоров С.К., Мирзоев Д.П.</i> (Республика Узбекистан) ИЗУЧЕНИЕ КОММУТАЦИОННЫХ УСТРОЙСТВ..... | 56 |
| <i>Kakhkhorov S.K., Jamilov Yu.Yu.</i> (Republic of Uzbekistan) OPPORTUNITIES OF THE FORMATION OF STUDENTS' COMPETENCE ON ALTERNATIVE ENERGY USING TRAINING SOFTWARE DEVICES / <i>Каххоров С.К., Жамилов Ю.Ю.</i> (Республика Узбекистан) ВОЗМОЖНОСТИ ФОРМИРОВАНИЯ КОМПЕТЕНТНОСТИ У СТУДЕНТОВ ПО АЛЬТЕРНАТИВНОЙ ЭНЕРГИИ С ИСПОЛЬЗОВАНИЕМ ПРОГРАММНЫХ СРЕДСТВ ОБУЧЕНИЯ..... | 61 |
| <i>Rasulova Z.D.</i> (Republic of Uzbekistan) DIDACTIC BASIS OF DEVELOPING CREATIVE THINKING OF FUTURE TEACHERS / <i>Расулова З.Д.</i> | |

| | |
|---|----|
| (Республика Узбекистан) ДИДАКТИЧЕСКИЕ ОСНОВЫ РАЗВИТИЯ У БУДУЩИХ УЧИТЕЛЕЙ КРЕАТИВНОГО МЫШЛЕНИЯ | 65 |
| <i>Ochilov Z.S., Hayitov O.A.</i> (Republic of Uzbekistan) INNOVATIVE FIELDS OF CREATIVE ACTIVITY OF PROFESSOR ADIBA SHARIPOVA / <i>Очиллов З.С., Хайитов О.А.</i> (Республика Узбекистан) ИННОВАЦИОННЫЕ СФЕРЫ ТВОРЧЕСКОЙ ДЕЯТЕЛЬНОСТИ ПРОФЕССОРА АДИБЫ ШАРИПОВОЙ | 69 |
| <i>Safarova D.S.</i> (Republic of Uzbekistan) PEDAGOGY OF COOPERATION AND EDUCATION DEVELOPMENT / <i>Сафарова Д.С.</i> (Республика Узбекистан) ПЕДАГОГИКА СОТРУДНИЧЕСТВА И РАЗВИТИЕ ОБРАЗОВАНИЯ | 71 |
| <i>Tukboeva D.Z.</i> (Republic of Uzbekistan) SOURCES OF FORMATION OF ECONOMIC CULTURE YOUNG PEOPLE IN THE WORKS OF EAST ENCYCLOPEDISTS SCIENTISTS / <i>Тукбоева Д.З.</i> (Республика Узбекистан) ИСТОКИ ФОРМИРОВАНИЯ ЭКОНОМИЧЕСКОЙ КУЛЬТУРЫ У МОЛОДЁЖИ В ТРУДАХ УЧЁНЫХ-ЭНЦИКЛОПЕДИСТОВ ВОСТОКА | 73 |
| <i>Ashrapov R.R.</i> (Republic of Uzbekistan) THE CULTURE OF BOOK READING IN THE FORMATION OF THE SOCIO-SPIRITUAL IMAGE OF YOUTH / <i>Ашрапов Р.Р.</i> (Республика Узбекистан) КУЛЬТУРА КНИГОЧТЕНИЯ В ФОРМИРОВАНИИ СОЦИАЛЬНО-ДУХОВНОГО ОБЛИКА МОЛОДЕЖИ | 75 |
| <i>Sharopova N.B.</i> (Republic of Uzbekistan) INTERACTIVE TECHNIQUES FOR TEACHING RUSSIAN LANGUAGE / <i>Шаропова Н.Б.</i> (Республика Узбекистан) ИНТЕРАКТИВНЫЕ ПРИЁМЫ ПРИ ОБУЧЕНИИ РУССКОМУ ЯЗЫКУ | 77 |
| <i>Yusupova I.B.</i> (Republic of Uzbekistan) SELF-KNOWLEDGE AND SELF-APPROVAL - KEY COMPONENTS OF THE MODERN PERSONALITY / <i>Юсупова И.Б.</i> (Республика Узбекистан) САМОПОЗНАНИЕ И САМОУТВЕРЖДЕНИЕ – КЛЮЧЕВЫЕ СОСТАВЛЯЮЩИЕ СОВРЕМЕННОЙ ЛИЧНОСТИ | 79 |
| <i>Jabborova D.F.</i> (Republic of Uzbekistan) INNOVATIVE TEACHING IMPROVEMENT TECHNOLOGIES / <i>Жабборова Д.Ф.</i> (Республика Узбекистан) ИННОВАЦИОННЫЕ ТЕХНОЛОГИИ СОВЕРШЕНСТВОВАНИЯ ОБУЧЕНИЯ | 81 |
| <i>Imotova G.F.</i> (Republic of Uzbekistan) LANGUAGE INTERACTION - AN IMPORTANT FACTOR FOR THE DEVELOPMENT OF PUPILS / <i>Имомова Г.Ф.</i> (Республика Узбекистан) ВЗАИМОДЕЙСТВИЕ ЯЗЫКОВ – ВАЖНЫЙ ФАКТОР РАЗВИТИЯ УЧЕНИКОВ | 83 |
| <i>Nematova N.K.</i> (Republic of Uzbekistan) MODERN TRENDS FOR FORMING ECONOMIC KNOWLEDGE IN A STUDENTING YOUTH / <i>Нематова Н.К.</i> (Республика Узбекистан) СОВРЕМЕННЫЕ ТЕНДЕНЦИИ ФОРМИРОВАНИЯ ЭКОНОМИЧЕСКИХ ЗНАНИЙ У УЧАЩЕЙСЯ МОЛОДЁЖИ | 85 |
| <i>Kurbonova M.A., Kurbonova N.A.</i> (Republic of Uzbekistan) POSSIBILITIES OF USING THE EDUCATIONAL COMPUTER PROGRAM IN MATHEMATICAL EDUCATION OF PRESCHOOLERS / <i>Курбонова М.А., Курбонова Н.А.</i> (Республика Узбекистан) ВОЗМОЖНОСТИ ИСПОЛЬЗОВАНИЯ УЧЕБНОЙ КОМПЬЮТЕРНОЙ ПРОГРАММЫ В МАТЕМАТИЧЕСКОМ ОБРАЗОВАНИИ ДОШКОЛЬНИКОВ | 87 |

ESSENTIAL SPECTRUM OF A 2×2 OPERATOR MATRIX AND THE FADDEEV EQUATION

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Abstract: we consider a 2×2 operator matrix \mathcal{A}_μ , $\mu > 0$ acting in the direct sum of one- and two-particle subspaces of a bosonic Fock space. It is related with the system of non conserved number of quasi-particles. We obtain an analogue of the Faddeev equation for the eigenfunctions of \mathcal{A}_μ . We describe the location of the essential spectrum of \mathcal{A}_μ via the spectrum of a family of generalized Friedrichs models. It is shown that the essential spectrum of \mathcal{A}_μ consists the union of at most 3 bounded closed intervals. We introduce new branches of the essential spectrum of \mathcal{A}_μ .

Keywords: operator matrix, bosonic Fock space, generalized Friedrichs model, essential spectrum, the Faddeev equation.

СУЩЕСТВЕННЫЙ СПЕКТР ОДНОЙ 2×2 ОПЕРАТОРНОЙ МАТРИЦЫ И УРАВНЕНИЕ ФАДДЕЕВА

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Аннотация: рассматривается 2×2 операторная матрица \mathcal{A}_μ , $\mu > 0$, действующая в прямой сумме одно- и двухчастичного подпространства бозонного пространства Фока. Оно связано с системой несохраняющихся чисел квазичастиц. Получен аналог уравнения Фаддеева для собственных функций оператора \mathcal{A}_μ . Местоположение существенного спектра оператора \mathcal{A}_μ описано с помощью спектра семейства обобщенных моделей Фридрихса. Показано, что существенный спектр оператора \mathcal{A}_μ состоит из объединения не более трех отрезков. Вводим новые ветви существенного спектра оператора \mathcal{A}_μ .

Ключевые слова: операторная матрица, бозонное пространство Фока, обобщенная модель Фридрихса, существенный спектр, уравнение Фаддеева.

It is well-known that every bounded linear operator acting in the direct sum of two Hilbert spaces always admits 2×2 block operator matrix representation [1]. The problems related with such matrices arise in statistical physics [2], solid-state physics [3] and the theory of quantum fields [4].

In the present paper we consider the family of 2×2 operator matrices \mathcal{A}_μ ($\mu > 0$ is a coupling constant) associated with the lattice systems describing two identical bosons and one particle,

another nature in interactions, without conservation of the number of particles. This operator acts in the direct sum of zero-, one- and two-particle subspaces of the bosonic Fock space and it is a lattice analogue of the spin-boson Hamiltonian. We derive an analogue of the Faddeev type system of integral equations for eigenvectors of \mathcal{A}_μ . We describe the location of the essential spectrum $\sigma_{ess}(\mathcal{A}_\mu)$ of \mathcal{A}_μ , via the spectrum of a family of generalized Friedrichs models $\mathcal{A}_\mu(k)$, $k \in T^d$. We introduce a new branches of $\sigma_{ess}(\mathcal{A}_\mu)$ and show that it consists the union of at most 3 bounded closed intervals. We find the discrete spectrum of \mathcal{A}_μ .

Let T^d be the d -dimensional torus, the cube $(-\pi, \pi]^d$ with appropriately identified sides equipped with its Haar measure. Let $L_2(T^d)$ be the Hilbert space of square integrable (complex) functions defined on T^d and $L_2^S((T^d)^2)$ be the Hilbert space of square integrable (complex) symmetric functions defined on $(T^d)^2$. Denote by H the direct sum of spaces $H_1 := L_2(T^d)$ and $H_2 := L_2^S((T^d)^2)$, that is, $H := H_1 \oplus H_2$. The spaces H_1 and H_2 are called one- and two-particle subspaces of a bosonic Fock space $F_S(L_2(T^d))$ over $L_2(T^d)$, respectively.

Let us consider a 2×2 operator matrices \mathcal{A}_μ acting in the Hilbert space H as

$$\mathcal{A}_\mu := \begin{pmatrix} A_{00} & \mu A_{01} \\ \mu A_{01}^* & A_{11} \end{pmatrix}$$

with the entries

$$(A_{11}f_1)(p) = w_1(k)f_1(p), \quad (A_{12}f_2)(p) = \int_{T^d} v(s)f_2(p, s)ds,$$

$$(A_{22}f_2)(p, q) = w_2(p, q)f_2(p, q), \quad f_i \in H_i, \quad i = 1, 2,$$

where $\mu > 0$ is a coupling constant, the functions $w_1(\cdot)$ and $w_2(\cdot; \cdot)$ are real-valued continuous functions on T^d and $(T^d)^2$ respectively. In addition the function $w_2(\cdot; \cdot)$ is a symmetric, that is, $w_2(p; q) = w_2(q; p)$ for any $p, q \in T^d$.

Under these assumptions the operator \mathcal{A}_μ is bounded and self-adjoint.

Let $H_0 := C$. To study the spectrum of the operator \mathcal{A}_μ we introduce a family of bounded self-adjoint operators (generalized Friedrichs models) $\mathcal{A}_\mu(k)$, $k \in T^d$ which acts in $H_0 \oplus H_1$ operator matrices

$$\mathcal{A}_\mu(k) := \begin{pmatrix} A_{00}(k) & \mu A_{01} \\ \mu A_{01}^* & A_{11}(k) \end{pmatrix}$$

with matrix elements

$$A_{00}(k)f_0 = w_1(k)f_0, \quad (A_{01}f_1) = \int_{T^d} v(t)f_1(t)dt, \quad (A_{11}f_2)(p) = w_2(k, p)f_2(p).$$

According to the Weyl theorem, for the essential spectrum of the operator $\mathcal{A}_\mu(k)$, we have $\sigma_{ess}(\mathcal{A}_\mu(k)) = [m(k); M(k)]$, where the numbers $m(k)$ and $M(k)$ are defined by $m(k) := \min_{p \in T^d} w_2(k, p)$ and $M(k) := \max_{p \in T^d} w_2(k, p)$.

For any fixed $k \in T^d$ we define an analytic function $\Delta_\mu(k; \cdot)$ (the Fredholm determinant associated with the operator $\mathcal{A}_\mu(k)$) in $C \setminus [m(k); M(k)]$

$$\Delta_\mu(k; z) := w_1(k) - z - \frac{\mu^2}{2} \int_{T^d} \frac{v^2(t) dt}{w_2(k, t) - z}.$$

Set $m := \min_{p, q \in T^d} w_2(p, q)$, $M := \max_{p, q \in T^d} w_2(p, q)$ and $\Lambda_\mu := \bigcup_{k \in T^d} \sigma_{disc}(\mathcal{A}_\mu(k))$.

We introduce a operator $T_\mu(z)$ acting in H_1 as

$$(T_\mu(z)g)(p) = \frac{\mu^2 v(p)}{2\Delta_\mu(p, z)} \int_{T^d} \frac{v(t)g(t)dt}{w_2(p, t) - z}, \quad z \notin \Sigma_\mu := [m; M] \cup \Lambda_\mu.$$

The following theorem [5-15] is an analog of the well-known Faddeev's result for the operator \mathcal{A}_μ .

Theorem 1. The number $z \in C \setminus \Sigma_\mu$ is an eigenvalue of the operator \mathcal{A}_μ if and only if the number $\lambda = 1$ is an eigenvalue of the operator $T_\mu(z)$. Moreover the eigenvalues z and 1 have the same multiplicities.

We point out that the equation $T_\mu(z)\varphi = \varphi$ is an analogue of the Faddeev type integral equation for eigenfunctions of the operator \mathcal{A}_μ .

Now we describe [11, 12, 16, 17] the location of the essential spectrum of the operator \mathcal{A}_μ by the spectrum of the family $\mathcal{A}_\mu(k)$ of generalized Friedrichs models.

Theorem 2. For the essential spectrum of \mathcal{A}_μ the equality $\sigma_{ess}(\mathcal{A}_\mu) = \Sigma_\mu$ holds. Moreover the set Λ_μ consists no more than three bounded closed intervals.

Following we introduce the new subsets of the essential spectrum of \mathcal{A}_μ .

Definition 1. The sets Λ_μ and $[m; M]$ are called two- and three-particle branches of the essential spectrum of \mathcal{A}_μ , respectively.

The definition of the set Λ_μ and the equality $\bigcup_{k \in T^d} [m(k); M(k)] = [m; M]$ together with Theorem 1 give the equality

$$\sigma_{ess}(\mathcal{A}_\mu) = \bigcup_{k \in T^d} \sigma(\mathcal{A}_\mu(k)). \quad (1)$$

Here the family of operators $\mathcal{A}_\mu(k)$ have a simpler structure than the operator \mathcal{A}_μ . Hence, in many instance, (1) provides an effective tool for the description of the essential spectrum. The spectral properties related with the threshold analysis of a family of 2×2 operator matrices were studied in [18-23]. In the paper [24] spectral inclusion property for diagonally dominant $n \times n$ unbounded operator matrices was studied.

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