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«AMALIY MATEMATIKA VA AXBOROT TEKNOLOGIYALARINING ZAMONAVIY MUAMMOLARI»
XALQARO ILMIY-AMALIY ANJUMAN

The poster features a blue background with several logos at the top right: the seal of the Republic of Uzbekistan, the seal of Tashkent State Transport University, the logo of Buxoro State University, and the seal of the Tashkent Mathematical Institute. The main title is centered in large, bold, dark blue font: «АМАЛИЙ МАТЕМАТИКА ВА АХБОРОТ ТЕХНОЛОГИЯЛАРИНИНГ ЗАМОНАВИЙ МУАММОЛАРИ». Below it, the subtitle «ХАЛҚАРО ИЛМИЙ-АМАЛИЙ АНЖУМАН» and the section title «МАТЕРИАЛЛАРИ» are also in large, bold, dark blue font. At the bottom left, the date «2022 йил, 11-12 май» is given. The bottom half of the poster shows a photograph of the modern white building of Buxoro State University with its name in blue letters on the facade. The overall design is professional and academic.

BUXORO – 2022

**ЎЗБЕКИСТОН РЕСПУБЛИКАСИ
ОЛИЙ ВА ЎРТА МАХСУС ТАЪЛИМ ВАЗИРЛИГИ
ЎЗБЕКИСТОН РЕСПУБЛИКАСИ ФАНЛАР АКАДЕМИЯСИ
В.И. РОМАНОВСКИЙ НОМИДАГИ МАТЕМАТИКА ИНСТИТУТИ
ЎЗБЕКИСТОН МИЛЛИЙ УНИВЕРСИТЕТИ
ТОШКЕНТ ДАВЛАТ ТРАНСПОРТ УНИВЕРСИТЕТИ
БУХОРО ДАВЛАТ УНИВЕРСИТЕТИ**

*Бухоро фарзанди, Беруний номидаги Давлат мукофоти лауреати, кўплаб
ёши изланувчиларнинг ўз йўлини топиб олишида раҳнамолик қилган етук
олим, физика-математика фанлари доктори Ғайбулла Назруллаевич
Салиховнинг 90 йиллик юбилейларига бағишиланади*

**АМАЛИЙ МАТЕМАТИКА ВА
АҲБОРОТ ТЕХНОЛОГИЯЛАРИНИНГ
ЗАМОНАВИЙ МУАММОЛАРИ**

**ХАЛҚАРО ИЛМИЙ-АМАЛИЙ АНЖУМАН
МАТЕРИАЛЛАРИ**

2022 йил, 11-12 май

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10. *J. E. Nápoles V., J. M. Rodríguez, J. M. Sigarreta*, New Hermite-Hadamard Type Inequalities Involving Non-Conformable Integral Operators, *Symmetry* 2019, 11, 1108; doi:10.3390/sym11091108

11. *M. Vivas-Cortez, O. J. Larreal B., J. E. Nápoles V.*, EXTREMAL SOLUTION TO GENERALIZED DIFFERENTIAL EQUATIONS UNDER INTEGRAL BOUNDARY CONDITION, *Journal of Mathematical Control Science and Applications*, Vol. 7 No. 1 (January-June 2021), 47-56

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QUARTIC NUMERICAL RANGE OF A TRIDIAGONAL 4×4 OPERATOR MATRICES

Latipov H.M., Rasulov T.H.

Bukhara State University

A block operator matrix is a matrix the entries of which are linear operators [1]. If the Hilbert space H is the product of four Hilbert spaces H_1, H_2, H_3 and H_4 , that is, $H = H_1 \oplus H_2 \oplus H_3 \oplus H_4$, then every bounded linear operator $A \in L(H)$ has a block operator matrix representation

$$A := \begin{pmatrix} A_{11} & A_{12} & A_{13} & A_{14} \\ A_{21} & A_{22} & A_{23} & A_{24} \\ A_{31} & A_{32} & A_{33} & A_{34} \\ A_{41} & A_{42} & A_{43} & A_{44} \end{pmatrix} \quad (1)$$

with bounded linear operators $A_{ij} \in L(H_j, H_i)$, $i, j = \overline{1, 4}$.

The operator A is a self-adjoint if and only if $A_{ij}^* = A_{ji}$ for all $i, j = \overline{1, 4}$. The following generalization of the numerical range of A takes into account the block structure (1) of A with respect to the decomposition $H = H_1 \oplus H_2 \oplus H_3 \oplus H_4$.

Let $S_{H_k} := \{x \in H_k : \|x\| = 1\}$, $k = \overline{1, 4}$ be the unit sphere in H_k and

$$S_H = S_{H_1} \times S_{H_2} \times S_{H_3} \times S_{H_4}.$$

Definition. For $f = (f_1, f_2, f_3, f_4) \in S_H$ we define 4×4 matrix

$$A_f := \begin{pmatrix} (A_{11}f_1, f_1) & (A_{12}f_2, f_1) & (A_{13}f_3, f_1) & (A_{14}f_4, f_1) \\ (A_{21}f_1, f_2) & (A_{22}f_2, f_2) & (A_{23}f_3, f_2) & (A_{24}f_4, f_2) \\ (A_{31}f_1, f_3) & (A_{32}f_2, f_3) & (A_{33}f_3, f_3) & (A_{34}f_4, f_3) \\ (A_{41}f_1, f_4) & (A_{42}f_2, f_4) & (A_{43}f_3, f_4) & (A_{44}f_4, f_4) \end{pmatrix} \in M_4(\mathbb{C}).$$

Then the set

$$\mathcal{W}^4(A) := \bigcup_{f \in S_H} \sigma_p(A_f)$$

is called the quartic numerical range of A (with respect to the block operator matrix representation (1)).

The block numerical range for $n \times n$ operator matrices was introduced in [2] for bounded entries and in [3] for unbounded entries.

In this note we consider the case where $A_{ij} = 0$ if $|i - j| \neq 1$ and $A_{ij}^* = A_{ji}$ if $|i - j| = 1$ for $i, j = \overline{1, 4}$. Our main results include a new formula for $\mathcal{W}^4(A)$ and an estimate for the bounds of A in terms of the quartic numerical range.

Let us introduce the following notations:

$$P(f) := |(A_{12}f_2, f_1)|^2 + |(A_{23}f_3, f_2)|^2 + |(A_{34}f_4, f_3)|^2,$$

$$Q(f) := |(A_{12}f_2, f_1)|^2 + |(A_{34}f_4, f_3)|^2;$$

$$E_1(f) := -\frac{\sqrt{2}}{2} \sqrt{P(f) + \sqrt{(P(f))^2 - 4Q(f)}}; \quad E_2(f) := -\frac{\sqrt{2}}{2} \sqrt{P(f) - \sqrt{(P(f))^2 - 4Q(f)}};$$

$$E_3(f) := \frac{\sqrt{2}}{2} \sqrt{P(f) - \sqrt{(P(f))^2 - 4Q(f)}};$$

$$E_4(f) := \frac{\sqrt{2}}{2} \sqrt{P(f) + \sqrt{(P(f))^2 - 4Q(f)}}.$$

The main result of this note is the following theorem.

Theorem. For the quartic numerical range $\mathcal{W}^4(A)$ of A we have

$$\mathcal{W}^4(A) = \bigcup_{k=1}^4 \bigcup_{f \in S_H} \{E_k(f)\}.$$

Moreover, for the lower and upper bounds of A the following estimates are holds:

$$\min \sigma(A) \geq \inf \mathcal{W}^4(A) = \inf_{f \in S_H} \bigcup \{E_1(f)\};$$

$$\max \sigma(A) \leq \sup \mathcal{W}^4(A) = \sup_{f \in S_H} \bigcup \{E_4(f)\}.$$

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THE GENERALIZED FRACTIONAL DIFFERENTIAL EQUATION OF LAGUERRE TYPE

Luciano M. Lugo Motta Bittencourt

Universidad Nacional del Nordeste Facultad de Ciencias Exactas y Naturales y Agrimensura Corrientes Capital, 3400, Argentina, E-mail: lmlmb@yahoo.com.ar

Abstract: In this paper, we present the Generalized Differential Equation of Laguerre Type writing with the generalized N-derivative and its associated Generalized Differential Equation, and we solve that using the generalized Laplace Transform and the power series method.

Key words: Generalized derivatives and integral, Laguerre Equation, Associated Laguerre Equation.

THE SPECTRUM OF THE DISCRETE SCHRÖDINGER OPERATOR WITH TWO-RANK PERTURBATION

Madatova F.A.

National University of Uzbekistan, E-mail: fotimamadatova2@gmail.com

The discrete Schrödinger operators have attracted considerable attention for both combinatorial Laplacians and quantum graphs; for some recent summaries refer to [1], [2] and the references therein. Particularly, eigenvalue behavior of discrete Schrödinger operators with small rank potentials are discussed in [3], [4], [5].

The one-particle discrete Schrödinger operator $h_{\lambda\mu}$ in the momentum representation is defined by

$$h_{\lambda\mu} = h_0 - V_{\lambda\mu},$$

where the non-perturbed operator h_0 acts on $L^2(\mathbb{T})$ as multiplication operator by the function $e(p)$:

$$(h_0 f)(p) = e(p)f(p), \quad f \in L^2(\mathbb{T}), \quad p \in \mathbb{T},$$

where

$$e(p) = 1 - \cos p, \quad p \in \mathbb{T}, \quad \mathbb{T} = (-\pi, \pi].$$

The potential operator of the form

$$(V_{\lambda\mu} f)(p) = \frac{\mu}{2\pi} \int_{\mathbb{T}} f(q) dq + \frac{\lambda}{2\pi} \int_{\mathbb{T}} e^{ix_0(p-q)} f(q) dq, \quad f \in L^2(\mathbb{T}), \quad p \in \mathbb{T}, \quad x_0 \in \mathbb{Z}.$$

For any $\lambda, \mu \in \mathbb{C}$, we define Fredholm determinant as a analytic function in $z \in \mathbb{C} \setminus [e_{\min}, e_{\max}]$ as

$$H_z(\lambda, \mu) = (\lambda - \gamma(z))(\mu - \gamma(z)) - \xi(z)$$

where

$$\gamma(z) = \frac{a(z)}{a^2(z) - b^2(z)}, \quad \xi(z) = \frac{b^2(z)}{(a^2(z) - b^2(z))^2}.$$

and

$$a(z) = \frac{1}{2\pi} \int_{\mathbb{T}} \frac{1}{e(q) - z} dq, \quad b(z) = \frac{1}{2\pi} \int_{\mathbb{T}} \frac{e^{ix_0 q}}{e(q) - z} dq.$$

Lemma. The number $z \in \mathbb{C} \setminus [e_{\min}, e_{\max}]$ is an eigenvalue of $h_{\lambda\mu}$ if only if $H(\lambda, \mu, z) = 0$.

We introduce the continuation of the function $H(\lambda, \mu, z)$ at the point $z = 0$ as follows

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